

Material Mechanics Research

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Final Report

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FOREWORD

This final technical report, entitled "Material Mechanic Research," presents the results of an in-house study performed under JON 2302M1G2 by AFRL/PRSM, Edwards AFB CA. The Principal Investigator/Project Manager for the Air Force Research Laboratory was Dr. C.T. Liu

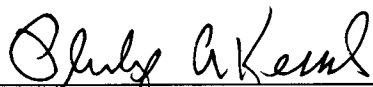
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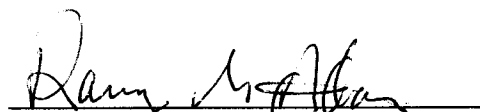
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GLOSSARY

psi	pounds per square inch
RL	Ravi-Lin
SFI	stress intensity factor
Tc	coefficient of thermal expansion
μ_o	undamaged material shear modulus
τ	relaxation time
κ_o	undamaged material bulk modulus

Introduction:

The goal of this program is to develop a basis for developing advanced crack growth and service life prediction technologies for predicting the service life of solid rocket motors. The objectives of this program are to : (1) gain a fundamental understanding of fracture and crack growth behavior in solid rocket motors; (2) investigate the effects of damage, material nonlinearity, pressure, and loading rate on crack growth behavior in a solid propellant; (3) simulate crack growth behavior and gain insight for improving crack growth resistance in solid propellants; and (4) determine the strain rate effect on the constitutive and fracture behavior of bi-material bond systems. The main issues in service life prediction of solid rocket motors are the lack of a fundamental understanding of crack growth behavior under service loading conditions and a reliable methodology to predict crack growth. The main technical challenges are microstructural effects on damage initiation and evolution, large and time dependent deformation, short crack and stress raiser interaction, and multi-layer structures with time-dependent material properties and property gradients. The program's basic approach involves a blend of analytical and experimental studies. In general, mechanisms and mechanics involved in cohesive fracture in a solid propellant and adhesive fracture in bond systems are emphasized. In this program, nonlinear viscoelasticity, fracture mechanics, experimental mechanics, damage mechanics, nondestructive testing and evaluation, and numerical modeling techniques will be used.

Objectives:

These research studies address a number of important subjects. There are four major tasks: (1) Task 1- investigating the effects of pressure, strain rate, and damage on crack growth behavior in a solid propellant, (2) Task 2 – developing a failure envelope to predict the critical stress for the onset of growth of short cracks, (3) Task 3 – investigating the three-dimensional effect on crack growth behavior, and (4) Task 4 - determining the strain rate effects on the constitutive behavior as well as on the local damage and deformation in bi-material, propellant/liner/propellant, bond systems. The results of these studies have the potential of becoming some of the most significant contributions to the rocket industry and research community.

Summary of Accomplishments:

Task 1: Investigating the Effects of Pressure, Strain Rate, and Damage on Crack growth Behavior in a Solid Propellant

In Task 1, experiments were conducted on uniaxial specimen with two different initial crack lengths (0.1 in. and 0.3 in.) at four different displacement rates (0.2 in/min, 2 in/min, 50 in/min, and 200 in/min) under three confined pressures (ambient, 500 psi, and 1000 psi). Experimental findings reveal that cracks grow slower under confined pressure. The decrease in crack growth rate under confined pressure is due to the suppression of the development of damage in the material and the decrease in strain gradient near the crack tip (1-6). The experimental findings also reveal that, after the crack propagates, the growth behavior of the two initial crack lengths considered are similar.

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In addition to conducting the experiments, the focus of research also centered on the effect of hydrostatic pressure on damage and the constitutive behavior of the material. Particularly, the study was conducted on modeling the effect of hydrostatic pressure on the progressive damage. A multi-scale approach was used along with a damage theory at the matrix material level. The multi-scale approach is based on the interconnection between the microlevel (i.e., particle and matrix material level) and the macrolevel (i.e., the particulate composite level). In order to model the effect of hydrostatic pressure on the damage, the damage theory previously developed was modified.

It has been shown that the damage function for solid propellants could be well expressed in terms of the total strain energy density. The strain energy has two parts: dilatational energy and distortional energy. While increasing distortional energy contributes to the damage, dilatational energy needs to be treated properly depending on positive and negative dilatation. If dilatation is positive, it is added to the distortional energy because it also causes damage. However, negative dilatation is not added to the distortional energy because it does not produce damage. Consequently, the strain energy damage function is expressed as :

$$f = u_d + H(e)u_v \quad (1)$$

where u_d and u_v are distortional and dilatational energy density, and H is the Heaviside unit step function defined such that $H(e) = 1$ when $e > 0$ and $H(e) = 0$ when $e < 0$. The parameter e denotes dilatation. Thus, hydrostatic pressure (negative dilatation) affects damage initiation and growth. The predicted stress-strain curves under hydrostatic pressure of 0 and 1000 psi, respectively, are plotted and compared to the experimental curves. The comparison is excellent so that the damage function of Equation (1) was proved to be valid.

The developed time-independent constitutive model was incorporated in a computer code, which was used to predict the initial crack length in high stress regions in an analog specimen under a multi-axial loading condition (7-10). The results are discussed in the following paragraphs.

Damage growth and crack initiation at a notch with stress concentration was studied using an analog specimen. For the analog specimen, there are two different notch radii, one large and one small. Therefore, there are two stress concentration locations in the specimen which are consistent with the results obtained from a photoelastic analysis. In the finite element analysis, 3-D solid elements with eight nodes per element were used. In the numerical analysis, a uniform displacement was applied to the slanted top boundary of the specimen and the symmetric boundary conditions were applied to the horizontal bottom boundary.

The notched specimen under tensile load with the displacement control was first studied without hydrostatic pressure. Before damage began in the specimen, the stress in the loading direction was the greatest at the element that was located at the intersection between the large and small radii of the notch. This element is called the ‘small radius element’ in the following discussion. Consequently, damage started to form at this site first. However, the stress level at the element on the symmetric boundary was quite comparable to that of the small radius element so that damage

followed subsequently on the symmetric boundary. The element on the symmetric boundary is called the 'large radius element' from now on. Eventually, damage saturated at the small radius element and the initial crack size, using the criterion discussed previously, was 1.7 mm long. The experimentally measured crack length was 1.86 mm in average so that there was a good agreement between the predicted and measured initial crack lengths.

The same kind of specimen was subjected to hydrostatic pressure of 1000 psi and also loaded in tension with the displacement control. For this case, the stress along the loading direction was the largest at the small radius element with initial damage. The large radius element competed with the small radius element in damage growth and the former had the damage saturation earlier than the latter. The initial crack size was 1.7 mm at section of the large radius element.

Task 2: Crack Instability and Growth Models

Sub-Task 1: Instability Criteria for Short Crack Growth

In this task, the instability criteria for the onset of growth of short cracks are developed. Based on experimental findings, we find that for short cracks, defined as having a crack length equal to or less than 0.1 in., classic fracture mechanics cannot be used to determine the critical condition for the onset of crack growth. In order to predict the onset of growth of short cracks, a linear fracture mechanics solution was modified to predict the behavior of short cracks. An effective crack length, which was equal to the original crack length plus the length of the failure process zone, was introduced into the Mode I stress intensity factor. By using the effective crack length, a reasonably good agreement exists between the fracture toughnesses for the onset of growth of short and long cracks. In addition, a failure envelope, based on theories of strength of material and fracture mechanics, was also developed. The developed failure envelope can be used to predict the critical stress for the onset of crack growth for short and long cracks under different strain rate and confining pressure conditions.

Sub-Task2: Determining the Damage Characteristics near the Crack Tip

In this task, Lockheed-Martin Research Laboratory's High-Resolution Digital Real-Time Radiographic System and acoustic imaging systems were used to determine the damage characteristics near the crack tip for different materials (11-13). The experimental results reveal that the critical damage intensity for the onset of crack growth is insensitive to loading history and material property. However, due to the viscoelastic nature of the materials, load history and time have a significant effect on the damage characteristics in the materials. Therefore, caution should be exercised when interpreting the nondestructive testing results.

Sub-Task 3: Predicting the Critical Inherent Initial Crack Length in the Material

In this task, a technique was developed (14), based on fracture mechanics and probabilistic mechanics, to predict the critical inherent initial crack size. The developed technique can be used

to predict the critical inherent initial crack size at different strain rates with good accuracy (15). For example, at 18.182 in./in./min strain rate, the predicted and the measured initial critical crack sizes are 0.146 in. and 0.132 in., respectively. The results also show that the critical inherent initial crack sizes are insensitive to strain rate. For example, at 0.067 in./in./min and 0.727 in./in./min strain rates the predicted critical inherent initial crack sizes are 0.13 in. and 0.119 in., respectively. In addition, the strain rate has no significant effect of the statistical distribution function of the critical inherent initial crack length. For the three strain rates considered, the statistical distribution functions of the critical inherent initial crack length follow the Second Asymptotic Distribution of the Maximum Value.

In addition, the results of the analysis reveal that the magnitudes of the confining pressure, ambient and 1000 psi, have no significant effect on the predicted critical initial crack length. It also reveals that the critical initial crack length and the statistical distribution function are insensitive to specimen's thickness, strain rate, and loading axially. Therefore, for the material investigated, the critical initial crack length can be considered a material property. The determination of the critical initial crack size and its statistical distribution function will make statistical and reliability analyses of crack growth feasible.

Sub-Task 4 Numerical Simulation of Crack Growth

In order to develop a fundamental understanding of the crack growth process, the crack-tip fields obtained from the numerical simulations were examined. The numerical simulation results suggest that the crack growth occurs in a discrete manner and upon the attainment of a critical state of strain at a critical distance ahead of the crack tip. This is consistent with the critical damage criterion that is used in the simulations.

In this sub-task, crack propagation at room temperature in a solid propellant was analyzed at different loading rates using a modified Ravi-Liu (RL) model by including time dependence of the homogenized material, i.e., particulate composite with a rubbery matrix and high stiffness particles. Experimental observations show that the volume dilatation Θ (damage parameter in RL model) (16) is independent of the loading rate and is purely a function of applied strain. In other words, volume dilatation is dependent only on mechanical deformation (strain). Taking these observations into account, the material time dependence is introduced through a simple Maxwell model for the undamaged material, and it is assumed that both the bulk and shear response of the material are governed by the same relaxation time, τ . The undamaged material's bulk (κ_0) and shear (μ_0) moduli are given in terms of its rubbery moduli, κ_∞ and μ_∞ and Prony coefficients κ_1 and μ_1 ,

$$\kappa_0(t) = \kappa_\infty + \kappa_1 \exp(-t/\tau) \quad \text{and} \quad \mu_0(t) = \mu_\infty + \mu_1 \exp(-t/\tau). \quad (2)$$

Stress components σ_{ij} are evaluated using the convolution integral,

$$\mathbf{s}_{ij}(t) = \int_0^t \mathbf{k}_o(t-t')g(\Theta)\mathbf{d}_{ij}dt' + \int_0^t 2\mathbf{m}_o(t-t')h(\Theta)\mathbf{e}_{ij}dt' \quad (3)$$

where g and h are the damage parameters, which are purely functions of the volume dilatation.

The modified viscoelastic RL model has been implemented in a special purpose displacement based finite element code, FEAP-SP. The damage parameters governing the degradation of the bulk and the shear moduli, g and h are determined using the uniaxial response and dilatation data as outlined in the original RL model. The finite element method uses the Newton-Raphson iterative procedure for achieving force equilibrium and Newton's iterative method for enforcing a plane stress condition. Constitutive update (stress, dilatation) is performed through integration of Equation (2) for a finite time step. Time steps are chosen to preserve accuracy in the constitutive update and to ensure quadratic convergence of the residual energy norm.

Crack propagation was simulated in specimens with edge cracks subjected to prescribed displacement rates at the boundaries. Due to the inherent nature of the singular stress fields near the crack tip, the loading rate varies by at least two orders of magnitude from the loading boundary to the crack tip. The relaxation and damage parameters determined from the uniaxial stress-strain curves are used in the finite element calculations to update the local moduli as well as damage. Crack initiation and propagation is performed using a node release technique. Attainment of critical dilatation is used as the crack initiation criteria. The FE model continues to check at the end of every time step elements that have attained the specified critical dilatation. Once the critical dilatation is reached in an element (usually at the crack tip), the reaction force is computed at the crack tip node. Then the computed nodal force replaces the crack tip displacement specified condition. Now, holding the far field displacement fixed at the time of crack initiation, the nodal force at the crack tip is relaxed over a specified number of steps (typically 20). At the end of this node release procedure, the crack is deemed to propagate across the element in which critical dilatation had been reached. The time taken to relax the nodal force is kept to a minimum, typically 0.1% of the time step, so as not to artificially relax the material properties. Once the node release has been accomplished, i.e., the nodal force reaching zero, the corresponding boundary condition for that node is changed to traction free. Then, the far field loading is continued as before until the critical dilatation is reached in an element and then the previously mentioned node release procedure is repeated. Hence, crack propagation in the finite element calculations is a discrete process and from the times at which the critical dilatation is reached at the crack tip positions, change in crack length as a function of time can be computed. At every load step, the reaction force is computed which would provide the boundary load as a function of time. Once the node release becomes untenable, i.e., attaining critical conditions in one or more elements without sustaining further loading, the crack propagation is said to be unstable and the computations are terminated.

Simulations of crack growth were carried out in a solid propellant at two different rates, 0.1 in/min and 0.5 in/min and room temperature (17)). Displacements were specified at the specimen boundaries normal to the crack. The load increases nearly linearly with increasing displacement with the crack blunting monotonically. Upon reaching the critical dilatation in the crack tip vicinity, crack growth is initiated. The initiation time is in fairly good agreement with experimentally measured values. Upon initiation, the load gradually level off and agrees closely with experimental observations. Upon crack initiation, the successive elements near the crack tip

fail and node release is used to propagate the crack. The crack exhibits blunting-sharpening-blunting sequence during crack growth indicating the discontinuous nature of crack propagation in solid propellants. The crack growth data from the numerical simulations indicate that the predictions show a similar trend for the experiments with higher crack propagation velocities. Crack tip fields have also been examined which show self-similarity in damage patterns (e.g., dilatation) during crack growth. These simulations have validated, for the first time, an approach for realistic modeling of crack growth and hence prediction of crack length as function of time and loading rates in solid rocket motors. Such a simulation capability could be a useful tool in predicting life of rocket motors in service.

Task 3: Photoelastic Analysis of Three-Dimensional Effects of Cracking of Motor Grain Geometries under Internal Pressure Loads

In order to obtain some insight into the three-dimensional effects of cracking of motor grain geometries under load, a series of experiments on photoelastic scale models of motor grain were conducted using the frozen stress method. The models were capped at the ends and pressurized internally above critical temperature after real cracks were introduced at fin tips in critical locations. After growing to desired size, pressure was reduced to stop growth and held through cooling.

Photoelastic analysis of the selected model geometry (18-19) show that there are two critical locations at a fin tip; one at the confluence of the edge radius of $R = 1.3$ mm with the radius $R = 11.0$ mm of the central part of the fin tip, and the other on the fin axis itself. There are two positions on each fin tip where the confluence of the two above noted radii exist. A crack emanating from such a position we call an off-axis crack. A crack at the other location on a fin axis we call a symmetric crack, which, in this case, is symmetric with respect to both load and geometry, and is a Class I crack which grows readily. Off-axis cracks, however, generally do not occupy principal planes of stress or planes of symmetry, and are generally called Class II cracks which must turn or kink to eliminate some Mode II before becoming purely Mode I at which time they will grow readily as Class I cracks. By placing both types of cracks in the same model separated by uncracked or plugged fins, (plugs are cylinders used to seal holes made to allow the insertion of a shaft which carries a blade to a critical locus at the fin tip), it is confirmed that the symmetric cracks penetrate to the outer surface of the model before the off-axis cracks have grown significantly. Data also reveal that a shear mode along with Mode I in the off-axis cracks before they are completely turned.

Experimental findings reveal that the turning effect in three dimensional off-axis cracks involves a shear mode except near the fin tip surface and, in non-brittle materials, turns rather than kinks sharply. Upon eliminating the shear mode, pure Mode I occurs all along the crack front in the Class I sense. It also reveals that cracks symmetric in both load and geometry will grow far more readily than off-axis cracks due to the shear mode effect in the latter even though the stress maybe higher at the off-axis locus in an uncracked fin.

If the above observations are projected to motor grain, then at sufficiently low temperatures, which stiffen the matrix to the level of the hard particles, the above-described behavior might occur. However, at higher temperatures, large microscopic shear mode effects would likely occur, during particulate rearrangement and dewetting, probably retarding crack growth.

Taken collectively, the above studies suggest the following observations: (1) Symmetric cracks are more dangerous than off-axis cracks even though they do not start at the locus of maximum stress in the uncracked model. (2) Off-axis cracks directed parallel to the fin axis are also dangerous but less so than the symmetric cracks for they will grow on slightly curved, and (3) While some of the cracks penetrated the outer wall in the depth direction, none of the cracks penetrated the length of the cylinder.

These results suggest that the practice of using a through-the-cylinder length crack in design maybe a substantial over design and suggests a comparison with deep semi-elliptic cracks as an alternative approach.

Task 4: Deformation and Fracture of Bonded Systems

Sub-Task 1: Determining the Stress Intensity Factors for Interfacial Cracks in Bi-Material Bonded Systems

The objectives of this study were to determine the three dimensional and residual stress effects on the distribution of the stress intensity factor (SIF) K_I along the interfacial cracks in bi-material specimens. In this task, both stress frozen and photoelastic techniques were used to determine the SIF along the interfacial cracks, through and part-through cracks, in bi-material specimens using a three-specimen test method. A three specimens test series was conducted for each crack depth involving (1) A homogeneous edge cracked specimen, (2) a homogeneous bonded edge cracked specimen, and (3) An edged cracked bimaterial specimen (19-22). After no load and fully loaded stress freezing cycles, slices were removed along the crack front. The results reveal that there is little variation, within experimental scatter, between the side slices and the center slice but center slice results are always higher so only center slice results are reported.

Experimental data indicate that the fringes are continuous across the upper bondline since the adhesive is chemically the same as the upper material. However, there is a disturbance along the lower bondline where the material mismatch occurs. This pattern contrasts from the results for through cracks. This led to significantly different values of K_I above and below the bondline for the part through crack. In order to determine which (if either) values of K_I was correct in this case, a separate experiment was performed on a model consisting of materials where there is no T_c mismatch and which had previously been found to yield good results for through cracked specimens. In the analysis, the K_I values on both sides of the crack were averaged. While no clear technical explanation can be made to support this step to correct asymmetry apparently due to T_c mismatch and perhaps other effects, nevertheless the test with no T_c mismatch mentioned above, showed reasonable agreement for this procedure.

Finally, when comparing results of the part through cracks with those for through cracks by matching a/t for the part through crack with a/w of the through crack at 0.25, the through crack yielded higher SIF values in the ratio of $1.82/1.21 = 1.50$ suggesting that the part through cracks are much less severe as expected.

Sub-Task 2 Deformation and Failure Mechanism of Propellant/Liner/Propellant Bonded Specimens

In this task, a series of experiments on propellant/liner/propellant bonded specimens were conducted at 0.01 in/min displacement rate. A computer aided speckle interferometry technique was used to determine the displacement and strain distributions in the specimen. Two interface debonding modes, debonding from the center and the corner of the interface of the specimen, are observed. These debonding modes appear to be related to the specimen geometry. In addition, the strain rates in the interphase and liner layers increase with increasing time, which are significant different from the constant applied strain rate (23).

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